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# Ultrasound velocity measurements in orbital-degenerate frustrated spinel $\text{MgV}_2\text{O}_4$

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**Abstract.** Ultrasound velocity measurements of the orbital-degenerate frustrated spinel  $\text{MgV}_2\text{O}_4$  are performed in the disorder-free high-purity single crystal which exhibits successive structural and antiferromagnetic phase transitions, and in the disorder-introduced single crystal which exhibits spin-glass-like behavior. The measurements reveal coexisting two types of anomalous temperature dependence of the elastic moduli in the cubic paramagnetic phase: Curie-type softening with decreasing temperature, and softening with a characteristic minimum with decreasing temperature. These elastic anomalies should respectively originate from the coexisting orbital fluctuations and spin-cluster excitations.

## 1. Introduction

Vanadate spinel  $AV_2O_4$  ( $A = \text{Zn}$  [1],  $\text{Mg}$  [2] and  $\text{Cd}$  [3]) is a geometrically frustrated magnet which undergoes a cubic-to-tetragonal structural transition at a temperature  $T_s$  and an antiferromagnetic (AF) transition at a low temperature  $T_N$ . For  $AV_2O_4$ , it is considered that the lowering of the lattice symmetry by the structural transition at  $T_s$  leads to the release of the frustration by the AF transition at  $T_N$  lower than  $T_s$ . Thus the interplay of spin, orbital, and lattice degrees of freedom should play a crucial role for the release of frustration in  $AV_2O_4$ . We are interested in the frustrated phase (the cubic paramagnetic (PM) phase) of  $AV_2O_4$ . And we performed ultrasound velocity measurements in the magnesium vanadate spinel  $\text{MgV}_2\text{O}_4$ . This compound exhibits successive phase transitions of a cubic-to-tetragonal structural transition at  $T_s = 65$  K and an AF transition at  $T_N = 42$  K [2]. For  $\text{MgV}_2\text{O}_4$ , it is known that a small amount of disorder suppresses the structural and magnetic phase transitions, and induces spin-glass-like behaviour at low temperatures [4]. In the present study, we performed the measurements in the disorder-free high-purity single crystal with the successive phase transitions at  $T_s = 65$  K and  $T_N = 42$  K, and in the disorder-introduced single crystal which exhibits the spin-glass-like behaviour below  $T_f = 12.5$  K.

## 2. Experimental

The ultrasound velocity measurements were performed in two different types of  $\text{MgV}_2\text{O}_4$  single crystals grown by the floating-zone method: the disorder-free high-purity single crystal with the successive structural and AF transitions at  $T_s = 65$  K and  $T_N = 42$  K, named here as “ordered  $\text{MgV}_2\text{O}_4$ ”,



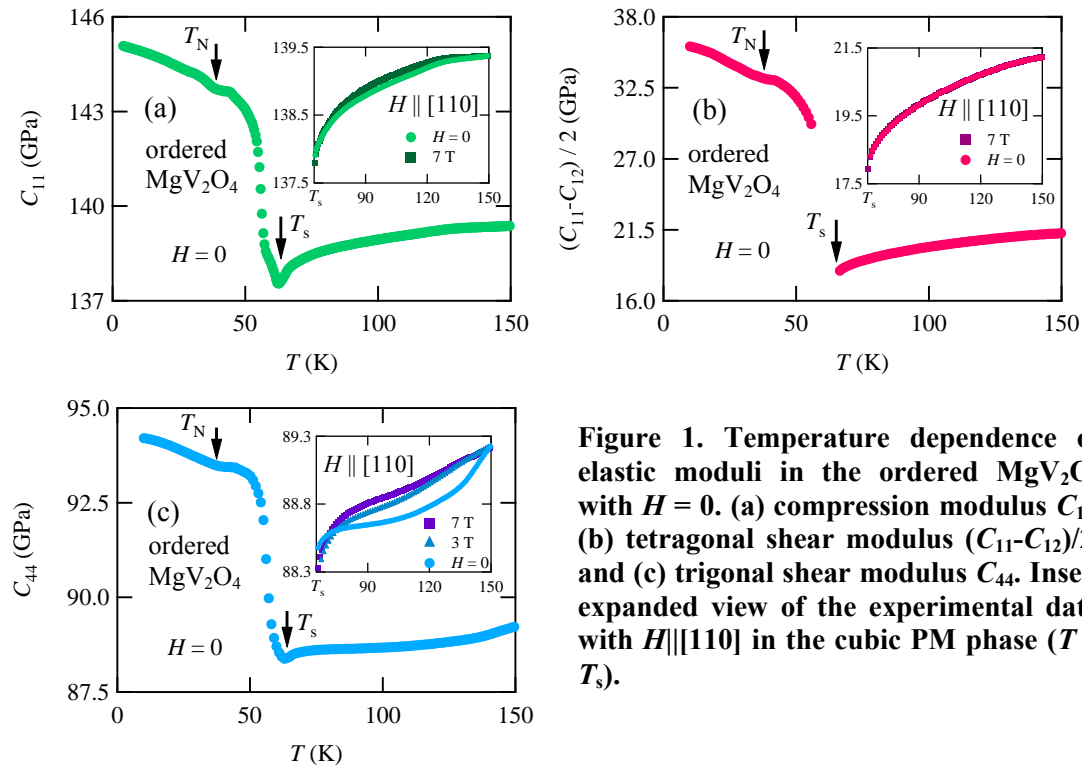
and the disorder-introduced single crystal with the spin-glass-like behavior below  $T_f = 12.5$  K in which  $\sim 3\%$  of V atoms in the octahedral sites are substituted by Mg atoms, named here as “disordered  $\text{MgV}_2\text{O}_4$ ” [4]. Temperature ( $T$ ) dependence of the ultrasound velocity was measured at  $T$  from 10 K to 150 K with magnetic field  $H \parallel [110]$  up to 7 T in all the symmetrically independent elastic moduli in the cubic crystal: compression modulus  $C_{11}$ , tetragonal shear modulus  $(C_{11}-C_{12})/2$ , and trigonal shear modulus  $C_{44}$ . The configuration of propagation  $\mathbf{k}$  and polarization  $\mathbf{u}$  of the sound wave for each elastic mode is summarized in Table 1.

**Table 1. Elastic modulus of cubic crystal and the corresponding configuration of sound wave with propagation  $\mathbf{k}$  and polarization  $\mathbf{u}$ .**

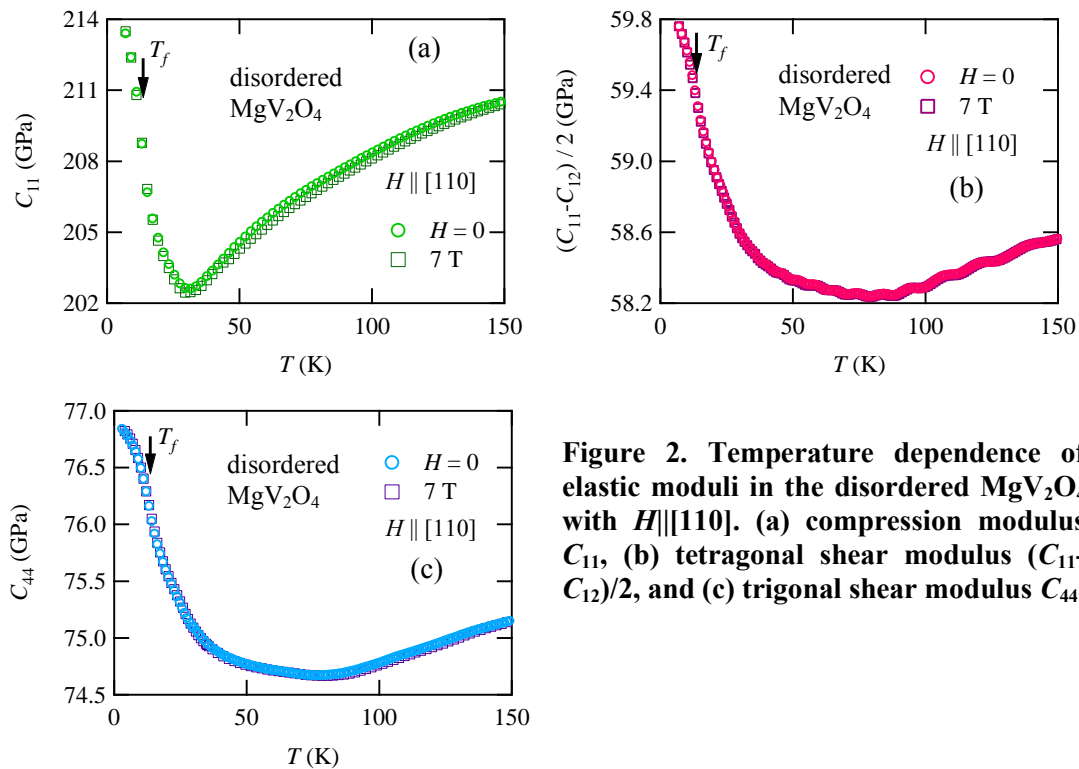
Elastic modulus	Sound wave	Propagation $\mathbf{k}$	Polarization $\mathbf{u}$
Compression modulus $C_{11}$	Longitudinal wave	[001]	[001]
Tetragonal shear modulus $(C_{11}-C_{12})/2$	Transverse wave	[1-10]	[110]
Trigonal shear modulus $C_{44}$	Transverse wave	[001]	[110]

### 3. Results and discussion

Figures 1(a), (b) and (c) respectively show  $T$  dependence of the compression modulus  $C_{11}$ , the tetragonal shear modulus  $(C_{11}-C_{12})/2$ , and the trigonal shear modulus  $C_{44}$  in the ordered  $\text{MgV}_2\text{O}_4$  with  $H = 0$ . All the elastic modes exhibit a jump at  $T_s$  and a discontinuous change at  $T_N$ , as marked by arrows in Figs. 1(a), (b), and (c). In the cubic PM phase ( $T > T_s$ ),  $C_{11}$  and  $(C_{11}-C_{12})/2$  exhibit huge Curie-type ( $\sim -1/T$ ) softening with decreasing  $T$ , which should be a precursor to the cubic-to-tetragonal lattice distortion at  $T_s$ . As shown in the insets to Figs. 1(a) and (b), this Curie-type softening is independent of  $H \parallel [110]$ , and thus should be driven by the coupling of lattice to orbital fluctuations which is hardly affected by the spin sector [5, 6, 7, 8]. On the other hand,  $C_{44}$  in the cubic PM phase ( $T > T_s$ ) exhibits non-monotonic softening with decreasing  $T$ : softening with convex curvature in  $\sim 80$  K  $< T < \sim 150$  K and Curie-type softening in  $T_s < T < \sim 80$  K. As shown in the inset to Fig. 1(c), the softening in  $C_{44}$  is sensitive to  $H \parallel [110]$ : the non-monotonic softening with concave curvature in the  $H$



**Figure 1. Temperature dependence of elastic moduli in the ordered  $\text{MgV}_2\text{O}_4$  with  $H = 0$ . (a) compression modulus  $C_{11}$ , (b) tetragonal shear modulus  $(C_{11}-C_{12})/2$ , and (c) trigonal shear modulus  $C_{44}$ . Inset: expanded view of the experimental data with  $H \parallel [110]$  in the cubic PM phase ( $T > T_s$ ).**



**Figure 2.** Temperature dependence of elastic moduli in the disordered  $\text{MgV}_2\text{O}_4$  with  $H \parallel [110]$ . (a) compression modulus  $C_{11}$ , (b) tetragonal shear modulus  $(C_{11}-C_{12})/2$ , and (c) trigonal shear modulus  $C_{44}$ .

$= 0$  data becomes closer to the Curie-type softening in the 7 T data. Thus, taking into account the  $H$  insensitivity of the Curie-type softening in  $C_{11}$  and  $(C_{11}-C_{12})/2$ , the non-monotonic softening in  $C_{44}$  should observe a superposition of  $H$ -sensitive concave  $T$  dependence and  $H$ -insensitive Curie-type softening. Furthermore, subtracting the component of the Curie-type softening from the observed non-monotonic softening in  $C_{44}$ , another component observed as the concave  $T$  dependence should be characterized as a softening with minimum with decreasing  $T$ . Such a softening-with-minimum anomaly is usually observed as a result of the coupling of lattice to magnetic excitations [9, 10, 11]. Taking into account that inelastic neutrons scattering experiments in  $AV_2O_4$  observed spin-cluster excitations in the cubic PM phase [1], the softening with minimum in  $C_{44}$  of the ordered  $\text{MgV}_2\text{O}_4$  should originate from the coupling of lattice to the spin-cluster excitations.

Figures 2(a), (b) and (c) respectively show  $T$  dependence of  $C_{11}$ ,  $(C_{11}-C_{12})/2$ , and  $C_{44}$  in the disordered  $\text{MgV}_2\text{O}_4$  with  $H \parallel [110]$  which behaves differently from that in the ordered  $\text{MgV}_2\text{O}_4$  shown in Figs. 1(a), (b), and (c). At  $T_f$  marked by arrows in Figs. 2(a), (b), and (c), all the elastic modes exhibit a small increase in slope. In the cubic PM phase ( $T > T_f$ ), more notably, all the elastic modes in the disordered  $\text{MgV}_2\text{O}_4$  exhibit the  $H$ -insensitive softening with minimum with decreasing  $T$  which should be driven by the coupling of lattice to the spin-cluster excitations.

We here note that the disordered  $\text{MgV}_2\text{O}_4$  exhibits the absence of the Curie-type softening, namely the absence of a precursor to structural transition, which is compatible with the absence of structural transition in the disordered  $\text{MgV}_2\text{O}_4$ . Taking into account the presence of the Curie-type softening in the ordered  $\text{MgV}_2\text{O}_4$  as shown in Fig. 1, the present study reveals that not only the structural transition but also its precursor (the Curie-type softening) in  $\text{MgV}_2\text{O}_4$  is sensitively suppressed by disorder. On the other hand, the softening with minimum is observed in both the ordered and the disordered  $\text{MgV}_2\text{O}_4$  indicating that the spin-cluster excitations are robust against disorder. Therefore the results in the present study strongly suggest the coexistence of the disorder-sensitive orbital fluctuations and the disorder-robust spin-cluster excitations in  $\text{MgV}_2\text{O}_4$ .

We note here that the component of the softening with minimum in  $C_{44}$  of the ordered  $\text{MgV}_2\text{O}_4$  seen in the inset to Fig. 1(c) is sensitive to  $H$ , whereas that in  $C_{11}$ ,  $(C_{11}-C_{12})/2$ , and  $C_{44}$  of the

disordered  $\text{MgV}_2\text{O}_4$  respectively seen in Figs. 2(a), (b), and (c) is insensitive to  $H$ . This difference in the  $H$ -sensitivity should also arise due to disorder, where the response of the spin-cluster-lattice coupling to  $H$  is quenched by the introduction of disorder. The detailed mechanism for this disorder effect remains to be elucidated. For instance, the excitation in  $\text{MgV}_2\text{O}_4$  might correctly be the orbital-spin-cluster excitation, and its change to the orbital-cluster excitation might occur by the introduction of disorder.

#### 4. Summary

We performed ultrasound velocity measurements of  $\text{MgV}_2\text{O}_4$  in the disorder-free high-purity single crystal which exhibits successive structural and AF phase transitions, and in the disorder-introduced single crystal which exhibits spin-glass-like behaviour. The measurements reveal coexisting two types of anomalous  $T$  dependence of the elastic moduli in the cubic PM phase: Curie-type softening with decreasing  $T$ , and softening with a characteristic minimum with decreasing  $T$ . These elastic anomalies should respectively originate from the coexisting disorder-sensitive orbital fluctuations and disorder-robust spin-cluster excitations.

#### 5. Acknowledgement

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#### References

- [1] Lee S H, Louca D, Ueda H, Park S, Sato T J, Isobe M, Ueda Y, Rosenkranz S, Zschack P, Iniguez J, Qiu Y, and Osborn R 2004 *Phys. Rev. Lett.* **93**, 156407.
- [2] Mamiya H, Onoda M, Furubayashi T, Tang J, and Nakatani I 1997 *J. Appl. Phys.* **81**, 5289.
- [3] Giovannetti G, Stroppa A, Picozzi S, Baldomir D, Pardo V, Blanco-Canosa S, Rivadulla F, Jodlauk S, Niermann D, Rohrkamp J, Lorenz T, Streltsov S, Khomskii D I, and Hemberger J 2011 *Phys. Rev. B* **83**, 060402(R).
- [4] Islam A T M N, Wheeler E M, Reehuis M, Siemensmeyer K, Tovar M, Klemke B, Kiefer K, Hill A H, and Lake B 2012 *Phys. Rev. B* **85**, 024203.
- [5] Kino Y, Luthi B, and Mullen M E 1972 *J. Phys. Soc. Jpn.* **33**, 687.
- [6] Kino Y, Luthi B, and Mullen M E 1973 *Solid State Commun.* **12**, 275.
- [7] Kataoka M and Kanamori J 1972 *J. Phys. Soc. Jpn.* **32**, 113.
- [8] Hazama H, Goto T, Nemoto Y, Tomioka Y, Asamitsu A, and Tokura Y 2000 *Phys. Rev. B* **62**, 15012.
- [9] Zherlitsyn S, Schmidt S, Wolf B, Schwenk H, Luthi B, Kageyama H, Onizuka K, Ueda Y, and Ueda K 2000 *Phys. Rev. B* **62**, R6097.